

Cubic Nonpolynomial Spline Approach to the Solution of a Second Order Two-Point Boundary Value Problem

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Abstract: Third and fourth order convergent methods based on cubic nonpolynomial spline function at midknotes are presented for the numerical solution of a second order two-point boundary value problem with Neumann conditions. Using this spline function a few consistency relations are derived for computing approximations to the solution of the problem. Convergence analysis of these methods is discussed two numerical examples are given to illustrate practical usefulness of the new methods. [Journal of American Science. 2010;6(12):297-302]. (ISSN: 1545-1003).

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1. Introduction:

In approximation theory spline functions occupy an important position having a number of applications, especially in the numerical solution of boundary-value problems. We shall consider a numerical solution of the following linear second order two-point boundary value problem, see [5].

$$y^{(2)} + f(x)y = g(x), \quad x \in [a, b] \quad (1.1)$$

Subject to Neumann boundary conditions:

$$y^{(1)}(a) - A_1 = y^{(1)}(b) - A_2 = 0 \quad (1.2)$$

Where A_i , $i = 1, 2$ are finite real constants. The functions $f(x)$ and $g(x)$ are continuous on the interval $[a, b]$. The analytical solution of (1.1) subjected to (1.2) cannot be obtained for arbitrary choices of $f(x)$ and $g(x)$.

The numerical analysis literature contains little on the solution of second order two-point boundary value problem (1.1) subjected to Neumann boundary conditions (1.2).while The linear second order two-point boundary value problem (1.1) subjected to Dirichlet boundary conditions solved by different types of spline functions, see [1, 7, 8, 9].

Ramadan et al. [5] solved the problem (1.1) subjected to (1.2) using quadratic polynomial spline, cubic polynomial spline and quadratic nonpolynomial spline at midknotes.

In this paper, we develop cubic nonpolynomial spline at midknotes to get smooth approximations for the solution of the problem (1.1) subjected to Neumann boundary conditions (1.2).

2. Derivation of the method:

We introduce a finite set of grid points x_i by dividing the interval $[a, b]$ into n equal parts.

$$x_i = a + ih, \quad i = 0, 1, \dots, n$$

$$x_0 = a, \quad x_n = b \quad \text{and} \quad h = \frac{b-a}{n} \quad (2.1)$$

Let $y(x)$ be the exact solution of the system (1.1) and (1.2) and S_i be an approximation to $y_i = y(x_i)$ obtained by the spline function $Q_i(x)$ passing through the points (x_i, s_i) and (x_{i+1}, s_{i+1}) .

Each nonpolynomial spline segment $Q_i(x)$ has the form.

$$Q_i(x) = a_i \sin k(x-x_i) + b_i \cos k(x-x_i) + c_i(x-x_i) + d_i, \quad i = 0, 1, \dots, n-1 \quad (2.2)$$

Where a_i , b_i , c_i and d_i are constants and k is the frequency of the trigonometric functions which will be used to raise the accuracy of the method and equation (2.2) reduces to cubic polynomial spline function in $[a, b]$ when $k \rightarrow 0$. Choosing the spline function in this form will enable us to generalize other existing methods by arbitrary choices of the parameters α and β which will be defined later at the end of this section. Thus, our cubic nonpolynomial spline is now defined by the relations:

$$(i) S(x) = Q_i(x), \quad x \in [x_i, x_{i+1}], \quad i = 0, 1, \dots, n-1$$

$$(ii) S(x) \in C^\infty[a, b] \quad (2.3)$$

The four coefficients in (2.2) need to be obtained in terms of

$S_{i+\frac{1}{2}}, D_i, M_{i+\frac{1}{2}}, T_i$ and T_{i+1} Where

$$\begin{aligned} (i) \quad Q_i(x_{i+\frac{1}{2}}) &= S_{i+\frac{1}{2}} \\ (ii) \quad Q_i^{(1)}(x_i) &= D_i \\ (iii) \quad Q_i^{(2)}(x_{i+\frac{1}{2}}) &= M_{i+\frac{1}{2}} \\ (iv) \quad Q_i^{(3)}(x_i) &= \frac{1}{2}[T_i + T_{i+1}] \end{aligned} \tag{2.4}$$

We obtain via a straightforward calculation

$$\begin{aligned} a_i &= \frac{-1}{2k^3}[T_{i+1} + T_i], \quad b_i = \frac{\tan \frac{\theta}{2}}{2k^3}[T_{i+1} + T_i] - \frac{\sec \frac{\theta}{2}}{k^2} M_{i+\frac{1}{2}} \\ c_i &= D_i + \frac{1}{2k^2}[T_{i+1} + T_i], \quad d_i = S_{i+\frac{1}{2}} + \frac{1}{k^2} M_{i+\frac{1}{2}} - \frac{h}{2} D_i - \frac{h}{4k^2}[T_{i+1} + T_i] \end{aligned} \tag{2.5}$$

Where $\theta = kh$ and $i = 0, 1, 2 \dots n-1$

Now using the continuity conditions (ii) and (2.3), that is the continuity of cubic nonpolynomial spline $S(x)$ and its first and second derivatives at the point (x_i, s_i) , where the two cubics $Q_{i-1}(x)$ and $Q_i(x)$ join, we can have

$$Q_{i-1}^{(m)}(x_i) = Q_i^{(m)}(x_i), \quad m = 0, 1, 2 \tag{2.6}$$

Using Eqs. (2.2), (2.4), (2.5) and (2.6) yield the relations:

$$\begin{aligned} \frac{h}{2}[D_i + D_{i-1}] &= (S_{i+\frac{1}{2}} - S_{i-\frac{1}{2}}) + \frac{1}{k^2} M_{i+\frac{1}{2}} [1 - \sec \frac{\theta}{2}] + \frac{1}{k^2} M_{i-\frac{1}{2}} [\cos \theta \sec \frac{\theta}{2} - 1] \\ &+ \left(\frac{\tan \frac{\theta}{2}}{2k^3} - \frac{h}{4k^2} \right) (T_{i-1} + 2T_i + T_{i+1}) \end{aligned} \tag{2.7}$$

$$\frac{h}{2}[D_i + D_{i-1}] = \frac{h \sin \theta \sec \frac{\theta}{2}}{2k} M_{i-\frac{1}{2}} \tag{2.8}$$

And

$$\frac{\tan \frac{\theta}{2}}{2k^3} (T_{i-1} + 2T_i + T_{i+1}) = \frac{\sec \frac{\theta}{2}}{k^2} M_{i+\frac{1}{2}} - \frac{\cos \theta \sec \frac{\theta}{2}}{k^2} M_{i-\frac{1}{2}} \tag{2.9}$$

From Eqs. (2.7) – (2.9) we get the following relation:

$$S_{i-\frac{3}{2}} - 2S_{i-\frac{1}{2}} + S_{i+\frac{1}{2}} = h^2 \left(\alpha M_{i-\frac{3}{2}} + \beta M_{i-\frac{1}{2}} + \alpha M_{i+\frac{1}{2}} \right), \quad i = 2, 3, \dots, n-1 \tag{2.10}$$

Where

$$\alpha = \frac{\theta - 2 \sin \frac{\theta}{2}}{2\theta^2 \sin \frac{\theta}{2}} \quad \text{And} \quad \beta = \frac{2\theta \sin^2 \frac{\theta}{2} + 4 \sin \frac{\theta}{2} - \theta(1 + \cos \theta)}{2\theta^2 \sin \frac{\theta}{2}}$$

And

$$M_i = -f_i S_i + g_i \quad \text{with} \quad f_i = f(x_i) \quad \text{and} \quad g_i = g(x_i)$$

The relation (2.10) gives $(n-2)$ linear algebraic equations in the (n) unknowns $S_{i+\frac{1}{2}}, i = 0, 1, 2, \dots, n-1$, so we need two more equations, one at each end of the range of integration for direct computation of $S_{i+\frac{1}{2}}$. These two equations are deduced by Taylor series and the method of undetermined coefficients. These equations are

$$-h S_0^{(1)} - S_{\frac{1}{2}} + S_{\frac{3}{2}} = h^2 \left(w_0 M_{\frac{1}{2}} + w_1 M_{\frac{3}{2}} + w_2 M_{\frac{5}{2}} + w_3 M_{\frac{7}{2}} \right), \quad \alpha i = 1 \tag{2.11}$$

And

$$S_{n-\frac{3}{2}} - S_{n-\frac{1}{2}} + h S_n^{(1)} = h^2 \left(w_0 M_{n-\frac{1}{2}} + w_1 M_{n-\frac{3}{2}} + w_2 M_{n-\frac{5}{2}} + w_3 M_{n-\frac{7}{2}} \right), \quad \alpha i = n \tag{2.12}$$

Where w_i 's will be determined later to get the required order of accuracy.

The local truncation errors $t_i, i = 1, 2, \dots, n$ associated with the scheme (2.10) – (2.12) can be obtained as follows, we rewrite the scheme (2.10) – (2.12) in the form

$$-h y_0^{(1)} - y_{\frac{1}{2}} + y_{\frac{3}{2}} = h^2 \left(w_0 y_{\frac{1}{2}}^{(2)} + w_1 y_{\frac{3}{2}}^{(2)} + w_2 y_{\frac{5}{2}}^{(2)} + w_3 y_{\frac{7}{2}}^{(2)} \right) + t_1, \quad \alpha i = 1 \tag{2.13}$$

$$y_{i-\frac{3}{2}} - 2y_{i-\frac{1}{2}} + y_{i+\frac{1}{2}} = h^2 \left(\alpha y_{i-\frac{3}{2}}^{(2)} + \beta y_{i-\frac{1}{2}}^{(2)} + \alpha y_{i+\frac{1}{2}}^{(2)} \right) + t_i, \quad \alpha i = 2, 3, \dots, n-1 \tag{2.14}$$

And

$$y_{n-\frac{3}{2}} - y_{n-\frac{1}{2}} + h y_n^{(1)} = h^2 \left(w_0 y_{n-\frac{1}{2}}^{(2)} + w_1 y_{n-\frac{3}{2}}^{(2)} + w_2 y_{n-\frac{5}{2}}^{(2)} + w_3 y_{n-\frac{7}{2}}^{(2)} \right) + t_n, \quad \alpha i = n \tag{2.15}$$

The terms $y_{i-\frac{1}{2}}^{(2)}$ and $y_{i-\frac{1}{2}}^{(2)} \dots$ in Eq. (2.14) are expanded around the point x_i using Taylor series and the expressions for $t_i, i = 2, \dots, n-1$ can be obtained. Also, expressions for $t_i; i = 1, n$ are obtained by expanding Eqs. (2.13) and (2.15) around

Then

$$\|E\|_{\infty} \leq \frac{\|M_0^{-1}\|_{\infty} \|T\|_{\infty}}{1 - \|M_0^{-1}\|_{\infty} \|J_0 + h^2 BF\|_{\infty}} \cong O(h^4) \tag{4.12}$$

We summarize the above results in the next theorem.

Theorem 4.1

Let $y(x)$ is the exact solution of the continuous boundary value problem (1.1) with the boundary condition (1.2) and

let $y_{i+1/2}, i = 0, 1, \dots, n - 1$, satisfies the discrete

boundary value problem (ii) in (3.1). Further, if

$$e_{i+1/2} = y_{i+1/2} - S_{i+1/2} \quad \text{then}$$

1- $\|E\|_{\infty} \cong O(h^3)$, for third order convergent method

2- $\|E\|_{\infty} \cong O(h^4)$, for fourth order convergent method

Which are given by (4.11) and (4.12), neglecting all errors due to round off.

5. Numerical examples and discussion:

We now consider two numerical examples illustrating the comparative performance of cubic nonpolynomial spline method (ii) in (3.1) over

quadratic nonpolynomial spline method and the two polynomial spline methods (quadratic and cubic). All calculations are implemented by MATLAB 7

Example 1

Consider the boundary value problem, see [5]

$$y^{(2)} + y = -1 \tag{5.1}$$

$$y^{(1)}(0) = \frac{1 - \cos(1)}{\sin(1)} = -y^{(1)}(1)$$

The analytical solution of (5.1) is

$$y(x) = \cos(x) + \frac{1 - \cos(1)}{\sin(1)} \sin(x) - 1 \tag{5.2}$$

Example 2

Consider the boundary value problem, see [5]

$$y^{(2)} + xy = (3 - x - x^2 + x^3) \sin(x) + 4x \cos(x) \tag{5.3}$$

$$y^{(1)}(0) = -1, y^{(1)}(1) = 2 \sin(1)$$

The analytical solution of (5.3) is

$$y(x) = (x^2 - 1) \sin(x) \tag{5.4}$$

The numerical results of examples 1 and 2 are presented in tables 1 and 2, respectively, for our fourth order method. A comparison between the method (2.10) and the existing methods in Ramadan et al. [5] are provided in tables 3 and 4.

Table 1: Approximate, Exact Solutions and Maximum errors (in absolute value) for Example 1 using our fourth order.

n	S_i (approximated)	y_i (Exact)	E (Error)
4	0.13068504600377	0.13060321651340	8.18295-5 ^a
8	0.13727099391989	0.13726907762415	1.91630-6
16	0.13893760135665	0.13893757908329	2.22734-8
32	0.00841355938534	0.00841356124929	1.86395-9
64	0.00423742716766	0.00423742736291	1.95255-10
128	0.00212635927486	0.00212635928910	1.42408-11

^a8.18295-5 = 8.18295*10⁻⁵

Table 2: Approximate, Exact Solutions and Maximum errors (in absolute value) for Example 2 using our fourth order.

n	S_i (approximated)	y_i (Exact)	E (Error)
4	- 0.35932989074946	- 0.35654365069809	2.78624-3
8	- 0.34264531263123	- 0.34258197359850	6.33390-5
16	- 0.29719707294621	- 0.29719640852255	6.64424-7
32	- 0.02582552960734	- 0.02582552960734	6.88566-8
64	- 0.01303052171412	- 0.01303052858255	6.86843-9
128	- 0.00654464520562	- 0.00654464571154	5.05929-10

Table 3: Maximum errors (in absolute value) for Example 1.

n	Our fourth order method	Our third order method	Quadratic nonpoly. [5]	Cubic polyn. [5]	Quadratic polyn. [5]
4	8.18295-5	8.18295-5	1.43181-3	2.85364-3	3.03488-3
8	1.91630-6	8.04854-6	1.75382-4	7.12633-4	7.69627-4
16	2.22734-8	5.91344-7	2.16003-5	1.78109-4	1.93094-4
32	1.86395-9	3.96416-8	2.67705-6	4.45241-5	4.83167-5
64	1.95255-10	2.56011-9	3.33110-7	1.11308-5	1.208186-5
128	1.42408-11	1.62945-10	4.15407-8	2.78270-6	3.02063-6

Table 4: Maximum errors (in absolute value) for Example 2.

n	Our fourth order method	Our third order method	Quadratic nonpoly. [5]	Cubic polyn. [5]	Quadratic polyn. [5]
4	2.78624-3	3.27323-3	2.2425-2	4.62182-2	4.94551-2
8	6.33390-5	3.03799-4	2.66946-3	1.15362-2	1.23088-2
16	6.64424-7	2.17464-5	3.24076-4	2.88302-3	3.08111-3
32	6.88566-8	1.43875-6	3.98761-5	7.20696-4	7.70391-4
64	6.86843-9	9.22972-8	4.94425-6	1.80171-4	1.92590-4
128	5.059294-10	5.84115-9	6.15517-7	4.50424-5	4.79946-5

6. Conclusion:

Two new methods are presented for solving second order two-point boundary value problem with Neumann conditions. These methods are shown to be optimal third and optimal fourth orders which are better than the two polynomial spline methods (quadratic and cubic splines) and quadratic nonpolynomial spline method. Moreover, nonpolynomial spline method has less computational cost over other polynomial spline methods. The obtained numerical results show that the proposed methods maintain a remarkable high accuracy which make them are very encouraging over other existing methods.

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